

# Lect 6: Quantum inverse method — Faddeev's method <sup>(1)</sup>

In last class, we have found

$$\left[ S_{j+1,j} S_{j+2,j} \cdots S_{N,j} S_{j,j} S_{2j} \cdots S_{j-1,j} \right] A_{\sigma_1 \cdots \sigma_N}^{(12 \cdots N, 12 \cdots N)} = e^{ik_j u} A_{\sigma_1 \cdots \sigma_N}^{(12 \cdots j, 12 \cdots j)}$$

with 
$$S_{ij} = \frac{k_i - k_j + ic P_{\sigma_i \sigma_j}}{k_i - k_j + ic}$$

In order to solve this problem, we use the following method of algebraic BA. We define an auxiliary space  $A$ , and  $\vec{c}$  as a Pauli matrix in such a space. Define  $P^{j,A} = \frac{1}{2}(1 + \vec{\sigma}_j \cdot \vec{c})$ , and the auxiliary  $S$ -matrix

$$S^{j,A}[u] = \frac{k_j - cu}{k_j - cu + ic} + \frac{ic}{k_j - cu + ic} P^{j,A}$$

The monodromy matrix

$$T(u) = S^{1,A}[u] S^{2,A}[u] \cdots S^{N,A}[u], \text{ where the matrix product only acts in the auxiliary space } A.$$

$$T_{\sigma_1 \sigma_2 \cdots \sigma_N u; \sigma'_1 \sigma'_2 \cdots \sigma'_N v} = S^{1,A}[u]_{\sigma_1 u; \sigma'_1 v} S^{2,A}[u]_{\sigma_2 u; \sigma'_2 v} \cdots S^{N,A}[u]_{\sigma_N u; \sigma'_N v}$$

Trace over the auxiliary space  $\Rightarrow$

$$\text{tr}_A [T(u_j)] = S_{j+1,j} S_{j+2,j} \cdots S_{N,j} S_{1,j} S_{2,j} \cdots S_{j-1,j}$$

where  $u_j = k_j / c$ .

Proof:  $\text{tr}_A [ T(u_j) ] = S_{\sigma_1 u_1, \sigma_1' u_2}^{1A} S_{\sigma_2 u_2, \sigma_2' u_3}^{2A} \dots S_{\sigma_N u_N, \sigma_N' u_1}^{NA}$

we can cyclic rotation over trace indices in the A-space

$$\text{tr}_A [ T(u_j) ] = \text{tr}_A [ S^{j+1,A} [u_j] S^{j+2,A} \dots S^{NA} S^{1A} \dots S^{j-1,A} S^{j,A} [u_j] ]$$

since  $S^{j,A} [u_j] = \frac{k_j - cu_j + ic p^{j,A}}{k_j - cu_j + ic} = p^{j,A}$

and we can prove an identity that  $S^{lA}(u_j) p^{jA} = p^{jA} S^{j,l}(u_j)$

Then  $\text{tr}_A [ T(u_j) ] = \text{tr}_A [ S^{j+1,A} [u_j] \dots S^{NA} [u_j] S^{1A} [u_j] \dots S^{j-1,A} [u_j] p^{jA} ]$   
 $= \text{tr}_A [ S^{j+1,A} [u_j] \dots S^{j-2,A} p^{jA} S^{j-1,j} [u_j] ]$   
 $= \text{tr}_A [ p^{jA} S^{j+1,j} [u_j] S^{j+2,j} [u_j] \dots S^{N,j} [u_j] S^{1,j} \dots S^{j-1,j} [u_j] ]$

$$p^{j,A} = \frac{1}{2} (1 + \vec{\sigma}_j \cdot \vec{c}) \Rightarrow \text{tr}_A [ p^{j,A} ] = 1$$

$\Rightarrow \text{tr}_A [ T(u_j) ] = S_{j+1,j} \dots S_{N,j} S_{1,j} \dots S_{j-1,j}$ , where  $u_j = k_j/c$ .

check:  $S^{lA}(u_j) p^{jA} = p^{jA} S^{j,l}(u_j) \rightarrow$  only need to check  $p^{lA} p^{jA} = p^{jA} p^{lA}$

$$p^{lA} p^{jA} = \frac{1}{4} (1 + \vec{\sigma}_l \cdot \vec{c})(1 + \vec{\sigma}_j \cdot \vec{c}) = \frac{1}{4} (1 + \vec{\sigma}_l \cdot \vec{c} + \vec{\sigma}_j \cdot \vec{c} + \sigma_l^\alpha \sigma_j^\beta c^\alpha c^\beta)$$

$$p^{jA} p^{lA} = \frac{1}{4} (1 + \vec{\sigma}_j \cdot \vec{c})(1 + \vec{\sigma}_l \cdot \vec{c}) = \frac{1}{4} (1 + \vec{\sigma}_j \cdot \vec{c} + \vec{\sigma}_l \cdot \vec{c} + \sigma_j^\alpha \sigma_l^\beta c^\alpha c^\beta)$$

$$c^\alpha c^\beta = \delta_{\alpha\beta} + i \epsilon_{\alpha\beta\gamma} c^\gamma, \quad \sigma_j^\alpha \sigma_j^\beta = \delta_{\alpha\beta} + i \epsilon_{\alpha\beta\gamma} \sigma_j^\gamma$$

$\Rightarrow$  it's easy to check

$$p^{lA} p^{jA} = p^{jA} p^{lA}$$

Then the Bethe ansatz equation is reduced to

$$\left[ \text{tr}_A T(k_j/c) \right]_{\sigma_1 \dots \sigma_N, \sigma'_1 \dots \sigma'_N} A_{\sigma'_1 \dots \sigma'_N} (12 \dots N; 12 \dots N) = e^{i k_j L} A_{\sigma_1 \dots \sigma_N} (12 \dots N, 12 \dots N)$$

We need to diagonalize  $\text{tr}_A T(k_j/c)$  simultaneously for  $j=1, 2, 3, \dots, N$ .

Define  $b(x) = \frac{-x}{-x+i}$ ,  $c(x) = \frac{i}{-x+i}$

and R-matrix in two auxiliary spaces  $A \otimes B$

$$R_{(u)}^{AB} = c(u) + b(u) P_{AB} = (b(u) + c(u) P_{AB}) P_{AB}$$

$$S^j_A(u) = b(u - k_j/c) + c(u - k_j/c) P_{jA}$$

$$S^j_B(u) = b(u - k_j/c) + c(u - k_j/c) P_{jB}$$

We have the fundamental commutation relation:

$$R^{AB}(u-v) [S^j_A(u) \otimes S^j_B(v)] = [S^j_A(v) \otimes S^j_B(u)] R^{AB}(u-v)$$

Proof: LHS =  $P_{AB} \left[ \frac{v-u+i P_{AB}}{v-u+i} \right] \left[ \frac{u_j-u+i P_{jA}}{u_j-u+i} \right] \left[ \frac{u_j-v+i P_{jB}}{u_j-v+i} \right]$

$$= \left[ \frac{v-u+i P_{AB}}{v-u+i} \right] \left[ \frac{u_j-u+i P_{jB}}{u_j-u+i} \right] P_{AB} \left[ \frac{u_j-v+i P_{jB}}{u_j-v+i} \right]$$

$$= \left[ \frac{v-u+i P_{AB}}{v-u+i} \right] \left[ \frac{u_j-u+i P_{jB}}{u_j-u+i} \right] \left[ \frac{u_j-v+i P_{jA}}{u_j-v+i} \right] P_{AB}$$

Compare with RHS  $\left[ \frac{u_j-v+i P_{jA}}{u_j-v+i} \right] \left[ \frac{u_j-u+i P_{jB}}{u_j-u+i} \right] \left[ \frac{v-u+i P_{AB}}{v-u+i} \right] P_{AB}$

check  $P_{AB} P_{Bj} P_{jA} = P_{AB} P_{jA} P_{AB} = P_{jA} P_{jB} P_{AB} \checkmark$

$$\begin{aligned}
 & (v-u) P_{Bj} P_{jA} + (u_j-u) P_{AB} P_{jA} + (u_j-v) P_{AB} P_{Bj} \\
 &= (v-u) P_{AB} P_{Bj} + (u_j-u) P_{AB} P_{jA} + (u_j-v) P_{AB} P_{Bj} \\
 &= (u_j-u) P_{AB} (P_{Bj} + P_{jA}) \\
 & (u_j-v) P_{jB} P_{AB} + (u_j-u) P_{jA} P_{AB} + (v-u) P_{jA} P_{jB} \leftarrow P_{jB} P_{AB} \\
 &= (u_j-v+v-u) P_{jB} P_{AB} + (u_j-u) P_{jA} P_{AB} \\
 &= (u_j-u) (P_{jB} + P_{jA}) P_{AB} = (u_j-u) P_{AB} (P_{jA} + P_{jB}) \checkmark
 \end{aligned}$$

terms only involving one permutation are the same  $\checkmark$

Now we can generalize it

$$\begin{aligned}
 & R_{(u-v)}^{AB} [ S^{j-1A}(u) S^{jA}(u) \otimes S^{j-1B}(v) S^{jB}(v) ] \\
 &= R_{(u-v)}^{AB} [ S^{j-1A}(u) \otimes S^{j-1B}(v), S^{jA}(u) \otimes S^{jB}(v) ] \\
 &= [ S^{j-1A}(v) \otimes S^{j-1B}(u) ] R^{AB}(u-v) [ S^{jA}(u) \otimes S^{jB}(v) ] \\
 &= [ S^{j-1A}(v) \otimes S^{j-1B}(u) ] [ S^{jA}(v) \otimes S^{jB}(u) ] R^{AB}(u-v) \\
 &= [ S^{j-1A}(v) S^{jA}(v) \otimes S^{j-1B}(u) S^{jB}(u) ] R^{AB}(u-v)
 \end{aligned}$$

$$\Rightarrow R^{AB}(u-v) [ S^{1A}(u) \dots S^{NA}(u) \otimes S^{1B}(v) \dots S^{NB}(v) ]$$

$$= [ S^{1A}(v) \dots S^{NA}(v) \otimes S^{1B}(u) \dots S^{NB}(u) ] R^{AB}(u-v)$$

RTT

i.e.  $R^{AB}(u-v) T^A(u) \otimes T^B(v) = T^A(v) \otimes T^B(u) R^{AB}(u-v)$

(5)

$R^{AB}$  is a  $4 \times 4$  matrix in the  $A \otimes B$  space

$$R^{AB}(u) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & C(u) & b(u) & 0 \\ 0 & b(u) & C(u) & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

and  $T^A(u)$  is a  $2 \times 2$  matrix in the  $A$  space

$$T(u) = \begin{pmatrix} A(u) & B(u) \\ C(u) & D(u) \end{pmatrix} \text{ and } A, B, C, D \text{ are } 2^N \times 2^N \text{ matrices in the physical Hilbert space}$$

Based on the "RTT" relation, we have quite a few usefully relations among operators  $A(u)$ ,  $B(u)$ ,  $C(u)$  and  $D(u)$ .

①  $A(u) A(v) = A(v) A(u)$

②  $A(u) B(v) = c(u-v) A(v) B(u) + b(u-v) B(v) A(u)$

③  $B(u) A(v) = c(u-v) B(v) A(u) + b(u-v) A(v) B(u)$

④  $B(u) B(v) = B(v) B(u)$

⑤  $A(v) B(u) = c(u-v) A(u) B(v) + b(u-v) B(u) A(v)$   
the same as ②

⑥  $C(u-v) [A(u) D(v) - A(v) D(u)] = b(u-v) [B(v) C(u) - C(u) B(v)]$

⑦  $C(u-v) [B(u) C(v) - B(v) C(u)] = b(u-v) [A(v) D(u) - D(u) A(v)]$

⑧  $B(v) D(u) = c(u-v) B(u) D(v) + b(u-v) D(u) B(v)$

$$(9) C(v)A(u) = c(u-v) C(u) A(v) + b(u-v) A(u) C(v)$$

$$(10) C(u-v) [C(u) B(v) - C(v) B(u)] = b(u-v) [D(v)A(u) - A(u)D(v)]$$

$$(11) C(u-v) [D(u)A(v) - D(v)A(u)] = b(u-v) [C(v) B(u) - B(u) C(v)]$$

$$(12) D(v) B(u) = c(u-v) D(u) B(v) + b(u-v) B(u) D(v)$$

$$(13) C(u) C(v) = C(v) C(u)$$

$$(14) C(u) D(v) = c(u-v) C(v) D(u) + b(u-v) D(v) C(u)$$

$$(15) D(u) C(v) = c(u-v) D(v) C(u) + b(u-v) C(v) D(u)$$

$$(16) D(u) D(v) = D(v) D(u)$$

The proof is straight forward from the RTT relation.

$$R^{AB}(u-v) T^A(u) \otimes T^B(v) = T^A(v) \otimes T^B(u) R^{AB}(u-v)$$

which is a 4x4 matrix identity in the  $A \otimes B$  space. Expand it and check term by term. See Appendix for proof. Now we will use it to prove

$$[\text{tr}_A T(u), \text{tr}_A T(v)] = 0 .$$

Proof:  $\text{tr}_A T(u) = A(u) + D(u)$ ,  $\text{tr}_A T(v) = A(v) + D(v)$

$$[\text{tr}_A T(u), \text{tr}_A T(v)] = [A(u) + D(u), A(v) + D(v)]$$

$$= [A(u), A(v)] + [D(u), D(v)] + (A(u)D(v) - A(v)D(u) + D(u)A(v) - D(v)A(u))$$

From relation (1)  $\Rightarrow [A(u), A(v)] = 0$

(1b)  $\Rightarrow [D(u), D(v)] = 0$

(6)  $\Rightarrow A(u)D(v) - A(v)D(u) = \frac{b(u-v)}{c(u-v)} [B(v)C(u) - C(u)B(v)]$

(11) exchanging  $u$  and  $v$

$D(v)A(u) - D(u)A(v) = \frac{b(v-u)}{c(v-u)} [C(u)B(v) - B(v)C(u)]$

since  $b(x)/c(x) = -x/i = ix \Rightarrow \frac{b(u-v)}{c(u-v)} = -\frac{b(v-u)}{c(v-u)}$

$\Rightarrow A(u)D(v) - A(v)D(u) = D(v)A(u) - D(u)A(v)$

$\Rightarrow [\text{tr}_A T(u), \text{tr}_A T(v)] = 0.$

★ In the A-space,  $S^j A(u) = b(u-u_j) + c(u-u_j) \frac{1}{2} (1 + \vec{\sigma}_j \cdot \vec{\tau}_A)$

$$= b(u-u_j) + \frac{c(u-u_j)}{2} \left[ (1 + \sigma_j^z) \frac{\tau_A^z + 1}{2} + (1 - \sigma_j^z) \frac{1 - \tau_A^z}{2} + 2\sigma_+ \frac{\tau_x - i\tau_y}{2} + 2\sigma_- \frac{\tau_x + i\tau_y}{2} \right]$$

$$S^j A(u) = b(u-u_j) + c(u-u_j) \begin{bmatrix} \frac{1 + \sigma_j^z}{2} & \sigma_- \\ \sigma_+ & \frac{1 - \sigma_j^z}{2} \end{bmatrix}$$

In the Physical Hilbert space,  $H_1 \otimes H_2 \dots \otimes H_n$ , we define the vacuum state  $|0\rangle = |\uparrow \uparrow \dots \uparrow\rangle$ .

Apply  $S^{jA}(u)$  on  $|0\rangle$ , we have

$$S^{jA}(u) |0\rangle = \begin{bmatrix} b(u-u_j) + c(u-u_j) & c(u-u_j) \sigma_- \\ c(u-u_j) \sigma_+ & b(u-u_j) \end{bmatrix} |0\rangle$$

$$= \begin{bmatrix} 1 & c(u-u_j) \sigma_-^j \\ 0 & b(u-u_j) \end{bmatrix} |0\rangle, \quad \begin{array}{l} \text{we have used} \\ b(u-u_j) + c(u-u_j) = 1 \end{array}$$

$$\Rightarrow T(u) |0\rangle = S^{1A}(u) S^{2A}(u) \dots S^{NA}(u) |0\rangle$$

$$= \begin{bmatrix} 1 & c(u-u_1) \sigma_-^1 \\ 0 & b(u-u_1) \end{bmatrix} \begin{bmatrix} 1 & c(u-u_2) \sigma_-^2 \\ 0 & b(u-u_2) \end{bmatrix} \dots \begin{bmatrix} 1 & c(u-u_N) \sigma_-^N \\ 0 & b(u-u_N) \end{bmatrix} |0\rangle$$

$$= \begin{pmatrix} 1 & B(u) \\ 0 & \prod_{j=1}^N b(u-u_j) \end{pmatrix} |0\rangle$$

*← complicated*  
*→ can be proved by induction.*

$$T(u) |0\rangle = \begin{pmatrix} 1 & B(u) \\ 0 & \prod_{j=1}^N b(u-u_j) \end{pmatrix} |0\rangle$$

Hence  $A(u) |0\rangle = |0\rangle$

$$D(u) |0\rangle = \prod_{j=1}^N b(u-u_j) |0\rangle$$

$$C(u) |0\rangle = 0$$

$$B(u) |0\rangle = \sum_j \# |\uparrow \dots \uparrow \downarrow \uparrow \dots \rangle \leftarrow \begin{array}{l} \text{one magnon} \\ \text{state.} \end{array}$$

*↑*  
*j<sup>th</sup>*



\*  $|0\rangle$  is the eigenstate of  $\text{tr}_A(T(u)) = A(u) + D(u)$   
with the eigenvalue  $1 + \prod_{j=1}^N b(u-u_j)$ .

\*  $B(u)$  operator behaves like  $S^-$  to create "magnon" excitation.

To get all the other eigenstates of  $\text{tr}_A(T(u))$ , we apply the flipping operator  $B$  on the "ferro" state:

$B(v_1) B(v_2) \dots B(v_M) |0\rangle$  is  $N-M$  spin-up,  $M$  spin-down.

We prove that it is an eigenstate of  $\text{tr} T(u)$ .

From relation (4)  $\Rightarrow B(u)B(v) = B(v)B(u) \leftarrow$  commute

$$\textcircled{8} \Rightarrow D(u)B(v) = \frac{1}{b(u-v)} B(v)D(u) - \frac{c(u-v)}{b(u-v)} B(u)D(v)$$

$\textcircled{3}$  exchange  $u \leftrightarrow v$

$$A(u)B(v) = \frac{1}{b(v-u)} B(v)A(u) - \frac{c(v-u)}{b(v-u)} B(u)A(v)$$

Check: the case that  $M=2$

$$A(u)B(v_1)B(v_2) = \left[ \frac{B(v_1)A(u)}{b(v_1-u)} - \frac{c(v_1-u)}{b(v_1-u)} B(u)A(v_1) \right] B(v_2)$$

$$= \left\{ \frac{1}{b(v_1-u)} B(v_1) \left[ \frac{1}{b(v_2-u)} B(v_2)A(u) - \frac{c(v_2-u)}{b(v_2-u)} B(u)A(v_2) \right] \right.$$

$$\left. - \frac{c(v_1-u)}{b(v_1-u)} B(u) \left[ \frac{1}{b(v_2-v_1)} B(v_2)A(v_1) - \frac{c(v_2-v_1)}{b(v_2-v_1)} B(v_1)A(v_2) \right] \right\}$$

$$D(u)B(v_1) \dots B(v_M) = \left[ \frac{1}{b(u-v_1)} B(v_1)D(u) - \frac{c(u-v_1)}{b(u-v_1)} B(u)D(v_1) \right] B(v_2)$$

$$= \left\{ \frac{1}{b(u-v_1)} B(v_1) \left[ \frac{1}{b(u-v_2)} B(v_2)D(u) - \frac{c(u-v_2)}{b(u-v_2)} B(u)D(v_2) \right] \right.$$

$$\left. - \frac{c(u-v_1)}{b(u-v_1)} B(u) \left[ \frac{1}{b(v_1-v_2)} B(v_2)D(v_1) - \frac{c(v_1-v_2)}{b(v_1-v_2)} B(v_1)D(v_2) \right] \right\}$$

$A(u)|0\rangle = |0\rangle$   
 $D(u)|0\rangle = \prod_{j=1}^N b(u-u_j)|0\rangle$

$$[A(u)+B(u)] B(v_1) B(v_2) |0\rangle = \left\{ \frac{1}{b(v_1-u)b(v_2-u)} + \frac{\prod_{j=1}^N b(u-u_j)}{b(u-v_1)b(u-v_2)} \right\} B(v_1) B(v_2) |0\rangle$$

+ unwanted terms.

If unwanted terms were zero, then  $B(v_1) \dots B(v_N) |0\rangle$  is  $\text{tr}_A T(u)$ 's eigenstate with the eigenvalue  $\frac{1}{b(v_1-u)b(v_2-u)} + \frac{\prod_{j=1}^N b(u-u_j)}{b(u-v_1)b(u-v_2)}$

Set  $u = u_j \Rightarrow \text{tr}_A T(u_j)$ 's eigenvalue  $\frac{1}{b(v_1-u_j)b(v_2-u_j)}$  since  $b(0)=0$

Then we have

$$e^{ik_j L} = \frac{1}{b(v_1-u_j)b(v_2-u_j)}$$

the 1st order unwanted terms

$$-\frac{c(v_1-u)}{b(v_1-u)} \left[ \frac{1}{b(v_2-v_1)} - \frac{\prod_{j=1}^N b(v_1-u_j)}{b(v_1-v_2)} \right] B(u) B(v_2) |0\rangle$$

Set  $\prod_{j=1}^N b(v_1-u_j) = \frac{b(v_1-v_2)}{b(v_2-v_1)} = \frac{\prod_{\alpha \neq 1} b(v_1-v_\alpha)}{\prod_{\alpha \neq 1} b(v_\alpha-v_1)} \Rightarrow$  this term vanish

check terms for  $B(v_1)B(u)|0\rangle$  — the calculation is complicated but since  $B(v_1)B(v_2) = B(v_2)B(v_1)$

we should arrive at

$$-\frac{c(v_2-u)}{b(v_2-u)} \left[ \frac{1}{b(v_1-v_2)} - \frac{\prod_{j=1}^N b(v_2-u_j)}{b(v_2-v_1)} \right] B(u) B(v_1) |0\rangle$$

(11)

$$\text{set } \prod_{j=1}^N b(v_2 - u_j) = \frac{b(v_2 - v_1)}{b(v_1 - v_2)} = \frac{\prod_{\alpha \neq 2} b(v_2 - v_\alpha)}{\prod_{\alpha \neq 2} b(v_\alpha - v_2)}$$

Then we can generalize to  $M$  - "magnon" case

$$\begin{aligned} & [A(u) + D(u)] B(v_1) B(v_2) \dots B(v_M) |0\rangle \\ &= \left\{ \prod_{\alpha=1}^M \frac{1}{b(v_\alpha - u)} + \frac{\prod_{j=1}^N b(u - u_j)}{\prod_{\alpha=1}^M b(u - v_\alpha)} \right\} B(v_1) \dots B(v_M) |0\rangle \\ & - \frac{C(v_1 - u)}{b(v_1 - u)} \left\{ \frac{1}{\prod_{\alpha=2}^M b(v_\alpha - v_1)} - \frac{\prod_{j=1}^N b(v_1 - u_j)}{\prod_{\alpha=2}^M b(v_1 - v_\alpha)} \right\} B(u) B(v_2) \dots B(v_M) |0\rangle \\ & - \frac{C(v_k - u)}{b(v_k - u)} \left\{ \frac{1}{\prod_{\alpha \neq k} b(v_\alpha - v_k)} - \frac{\prod_{j=1}^N b(v_k - u_j)}{\prod_{\alpha \neq k} b(v_k - v_\alpha)} \right\} B(v_1) \dots B(u) B(v_{k+1}) \dots B(v_M) |0\rangle \end{aligned}$$

$\leftarrow A(v_1)|0\rangle$        $\leftarrow D(v_1)|0\rangle$   
 $\leftarrow A(v_k)|0\rangle$        $\leftarrow D(v_k)|0\rangle$

Due to the permutation symmetry

Hence, we set

$$\prod_{\alpha=1}^M \frac{1}{b(v_\alpha - u_j)} = e^{ik_j L} \quad \text{for } j=1, 2, \dots, N$$

$$\prod_{j=1}^N b(v_k - u_j) = \frac{\prod_{\alpha \neq k} b(v_k - v_\alpha)}{\prod_{\alpha \neq k} b(v_\alpha - v_k)} \quad \text{for } k=1, \dots, M$$

All unwanted terms vanish, and

$$\text{tr}_A [T(u_j)] \left\{ B(v_1) B(v_2) \dots B(v_M) |0\rangle \right\} = e^{ik_j L} \left[ B(v_1) \dots B(v_M) |0\rangle \right]$$

Plug in  $u_j = \frac{k_j}{c}$ ,  $v_\alpha = \frac{\Lambda_\alpha}{c} + \frac{i}{2}$ ,  $b(u) = \frac{-u}{-u+i}$  (12)

$$\Rightarrow b(v_\alpha - u_j) = \frac{-\frac{\Lambda_\alpha}{c} + \frac{k_j}{c} - \frac{i}{2}}{-\frac{\Lambda_\alpha}{c} + \frac{k_j}{c} + \frac{i}{2}} = \frac{k_j - \Lambda_\alpha - \frac{i c}{2}}{k_j - \Lambda_\alpha + \frac{i c}{2}}$$

$$b(v_\alpha - v_k) = \frac{\Lambda_k - \Lambda_\alpha}{\Lambda_k - \Lambda_\alpha + i c}$$

$$b(v_k - v_\alpha) = \frac{\Lambda_\alpha - \Lambda_k}{\Lambda_\alpha - \Lambda_k + i c}$$

$$\Rightarrow e^{i k_j L} = \prod_{\alpha=1}^M \frac{k_j - \Lambda_\alpha + \frac{i c}{2}}{k_j - \Lambda_\alpha - \frac{i c}{2}} \quad \leftarrow \text{for each } k_j \text{ } j=1, \dots, N$$

$$\prod_{j=1}^N \frac{(k_j - \Lambda_k) - \frac{i c}{2}}{k_j - \Lambda_k + \frac{i c}{2}} = \prod_{\alpha=1, \alpha \neq k}^M \frac{\Lambda_\alpha - \Lambda_k - i c}{\Lambda_\alpha - \Lambda_k + i c}$$

or rewrite

$$\prod_{j=1}^N \frac{k_j - \Lambda_\alpha + \frac{i c}{2}}{k_j - \Lambda_\alpha - \frac{i c}{2}} = - \prod_{\beta=1}^M \frac{\Lambda_\beta - \Lambda_\alpha + i c}{\Lambda_\beta - \Lambda_\alpha - i c} \quad \text{for } \alpha=1, \dots, M$$

Appendix:

(13)

$$\left[ R^{AB}(u-v) T^A(u) T^B(v) \right]_{uw, u'w'} = \left[ c(u-v) \delta_{uw, u''w''} + b(u-v) P_{uw, u''w''} \right]$$

$$= c(u-v) T_{uu'}^A(u) T_{ww'}^B(v) + b(u-v) T_{\omega u'}^A(u) T_{u\omega'}^B(v)$$

$\begin{matrix} \diagup & \diagdown \\ T^A(u) & T^B(v) \\ u''u' & \omega''\omega' \end{matrix}$

$$\left[ T^A(v) T^B(u) R^{AB}(u-v) \right]_{uw, u'w'} = T_{uu''}^A(v) T_{\omega\omega''}^B(u) \left[ c(u-v) \delta_{u''w'', u'w'} + b(u-v) P_{u''w'', u'w'} \right]$$

$$= c(u-v) T_{u\omega'}^A(u) T_{\omega w'}^B(u) + b(u-v) T_{u\omega'}^A(v) T_{\omega u'}^B(u)$$

$uw, u'w'$   
 (11, 11):  $T_{11}(u) T_{11}(v) = T_{11}(v) T_{11}(u) \Rightarrow \underline{A(u) A(v) = A(v) A(u)}$ ,

(11, 12)  $c(u-v) T_{11}(u) T_{12}(v) + b(u-v) T_{11}(u) T_{12}(v)$   
 $= c(u-v) T_{11}(v) T_{12}(u) + b(u-v) T_{12}(v) T_{11}(u)$

$\Rightarrow \underline{A(u) B(v) = c(u-v) A(v) B(u) + b(u-v) B(v) A(u)}$

(11, 13)  ~~$c(u-v) T_{11}(u) T_{13}(v) + b(u-v) T_{11}(u) T_{13}(v)$   
 $= c(u-v) T_{11}(v) T_{13}(u) + b(u-v) T_{13}(v) T_{11}(u)$~~

$\Rightarrow c(u-v) T_{12}(u) T_{11}(v) + b(u-v) T_{12}(u) T_{11}(v)$

$= c(u-v) T_{12}(v) T_{11}(u) + b(u-v) T_{11}(v) T_{12}(u)$

$\Rightarrow \underline{B(u) A(v) = c(u-v) B(v) A(u) + b(u-v) A(v) B(u)}$

$\Rightarrow A(v) B(u) = \frac{1}{b(u-v)} B(u) A(v) - \frac{c(u-v)}{b(u-v)} B(v) A(u)$

uw u'w'  
(11, 22)

$$\begin{aligned}
 & c(u-v) T_{12}(u) T_{12}(v) + b(u-v) T_{12}(u) T_{12}(v) \\
 & = c(u-v) T_{12}(v) T_{12}(u) + b(u-v) T_{12}(v) T_{12}(u)
 \end{aligned}
 \left. \vphantom{\begin{aligned} & c(u-v) T_{12}(u) T_{12}(v) + b(u-v) T_{12}(u) T_{12}(v) \\ & = c(u-v) T_{12}(v) T_{12}(u) + b(u-v) T_{12}(v) T_{12}(u) \end{aligned}} \right\} \Rightarrow \boxed{\begin{aligned} & B(u) B(v) \\ & = B(v) B(u) \end{aligned}}$$

uw u'w'  
(12, 11)

$$\begin{aligned}
 & c(u-v) T_{11}(u) T_{21}(v) + b(u-v) T_{21}(u) T_{11}(v) \\
 & = c(u-v) T_{11}(v) T_{21}(u) + b(u-v) T_{11}(v) T_{21}(u)
 \end{aligned}$$

$$\underline{c(u-v) A(u) B(v) + b(u-v) B(u) A(v) = A(v) B(u)}$$

uw u'w'  
(12, 12)

$$\begin{aligned}
 & c(u-v) T_{11}(u) T_{22}(v) + b(u-v) T_{21}(u) T_{12}(v) \\
 & = c(u-v) T_{11}(v) T_{22}(u) + b(u-v) T_{12}(v) T_{21}(u)
 \end{aligned}$$

$$\star \underline{c(u-v) A(u) D(v) + b(u-v) C(u) B(v) = c(u-v) A(v) D(u) + b(u-v) B(v) C(u)}$$

uw u'w'  
12 21

$$\begin{aligned}
 & c(u-v) T_{12}(u) T_{21}(v) + b(u-v) T_{22}(u) T_{11}(v) \\
 & = c(u-v) T_{12}(v) T_{21}(u) + b(u-v) T_{11}(v) T_{22}(u)
 \end{aligned}$$

$$\star \underline{c(u-v) B(u) C(v) + b(u-v) D(u) A(v) = c(u-v) B(v) C(u) + b(u-v) A(v) D(u)}$$

uw, u'w'  
12 22

$$\begin{aligned}
 & c(u-v) T_{12}(u) T_{22}(v) + b(u-v) T_{22}(u) T_{12}(v) \\
 & = c(u-v) T_{12}(v) T_{22}(u) + b(u-v) T_{12}(v) T_{22}(u)
 \end{aligned}$$

$$\underline{c(u-v) B(u) D(v) + b(u-v) D(u) B(v)} \Rightarrow \underline{B(v) D(u)}$$

$$\begin{aligned}
 D(u) B(v) &= \frac{B(v) D(u)}{b(u-v)} \\
 &- \frac{c(u-v)}{b(u-v)} B(u) D(v)
 \end{aligned}$$

$u \quad w \quad u'w'$   
 $21 \quad 11$

$$c(u-v) T_{21}(u) T_{11}(v) + b(u-v) T_{11}(u) T_{21}(v)$$

$$= c(u-v) T_{21}(v) T_{11}(u) + b(u-v) T_{21}(v) T_{11}(u)$$

$$c(u-v) c(u) A(v) + b(u-v) A(u) c(v) = c(v) A(u)$$

---

21 12  $c(u-v) T_{21}(u) T_{12}(v) + b(u-v) T_{11}(u) T_{22}(v)$

$$= c(u-v) T_{21}(v) T_{12}(u) + b(u-v) T_{22}(v) T_{11}(u)$$

$$c(u-v) c(u) B(v) + b(u-v) A(u) D(v) = c(u-v) c(v) B(u) + b(u-v) D(v) A(u)$$

---

21 21  $c(u-v) T_{22}(u) T_{11}(v) + b(u-v) T_{12}(u) T_{21}(v)$

$$= c(u-v) T_{22}(v) T_{11}(u) + b(u-v) T_{21}(v) T_{12}(u)$$

$$c(u-v) D(u) A(v) + b(u-v) B(u) c(v) = c(u-v) D(v) A(u) + b(u-v) c(v) B(u)$$

---

21 22  $c(u-v) T_{22}(u) T_{12}(v) + b(u-v) T_{12}(u) T_{22}(v)$

$$= c(u-v) T_{22}(v) T_{12}(u) + b(u-v) T_{22}(v) T_{12}(u)$$

$$c(u-v) D(u) B(v) + b(u-v) B(u) D(v) = c(u-v) D(v) B(u) + b(u-v) D(v) B(u)$$

---

$$= D(v) B(u)$$

22 11  $c(u-v) T_{21}(u) T_{21}(v) + b(u-v) T_{21}(u) T_{21}(v)$

$$= c(u-v) T_{21}(v) T_{21}(u) + b(u-v) T_{21}(v) T_{21}(u)$$

$$c(u) c(v) = c(v) c(u)$$

$u \ w \ u' \ w'$

22, 12

$$c(u-v) T_{21}(u) T_{22}(v) + b(u-v) T_{21}(u) T_{22}(v)$$

$$= c(u-v) T_{21}(v) T_{22}(u) + b(u-v) T_{22}(v) T_{21}(u)$$

$$\Rightarrow \underline{c(u) D(v) = c(u-v) c(v) D(u) + b(u-v) D(v) c(u)}$$

$u \ w \ u' \ w'$

22 21

$$c(u-v) T_{22}(u) T_{21}(v) + b(u-v) T_{22}(u) T_{21}(v)$$

$$= c(u-v) T_{22}(v) T_{21}(u) + b(u-v) T_{21}(v) T_{22}(u)$$

$$\underline{D(u) c(v) = c(u-v) D(v) c(u) + b(u-v) c(v) D(u)}$$

22 22

$\Rightarrow$

$$\underline{D(u) D(v) = D(v) D(u)}$$



Next we prove :

$$S^{\pm} B(u_1) \dots B(u_m) |0\rangle,$$

$$S^3 B(u_1) \dots B(u_m) |0\rangle = \frac{N-2m}{2} B(u_1) \dots B(u_m) |0\rangle,$$

where  $S^{\pm} = S^1 \pm iS^2$ , i.e.  $B(u_1) \dots B(u_m)$  is  $S_{tot} = \frac{N-2m}{2}$  state.

Proof:  $[S^{j,A}(u), S_j^n] = [b(u-u_j) + c(u-u_j) \frac{1+\vec{z} \cdot \vec{\sigma}_j}{2}, \frac{\sigma_j^n}{2}]$  (n=1,2,3  
spin  
orientation)

$$= \frac{1}{4} c(u-u_j) \tau^m [\sigma_j^m \sigma_j^n] \leftarrow 2i \epsilon^{mnl} \tau^m \sigma^l$$

$$= -\frac{1}{4} c(u-u_j) \sigma_j^l [\tau_j^l \tau_j^n] \leftarrow 2i \epsilon^{lnm} \sigma^l \tau^m$$

$$= - [S^{j,A}(u), \frac{1}{2} \tau_j^n] = [S^{j,A}(u), S_j^n]$$

$$T(u) = S^{1A}(u) S^{2A}(u) \dots S^{NA}(u)$$

$$[T(u), S_{tot}^n] = \sum_{j=1}^N S^{jA} \dots [S^{jA}, S_{tot}^n] \dots S^{jA}$$

$$= - \sum_j L_j \dots [L_j, \frac{\tau_j^n}{2}] \dots L_N = - [T(u), \frac{1}{2} \tau^n]$$

in the A space:  $T(u)_{\alpha\beta} \leftarrow$  explicitly write the matrix element.

$$\Rightarrow [T(u)_{\alpha\beta}, S_{tot}^n] = - [T(u) \frac{1}{2} \tau^n] = \frac{1}{2} [ \tau_{\alpha\beta}^n T_{\beta'\beta} - T_{\alpha\beta'} \tau_{\beta'\beta}^n ]$$

$$T_N = \begin{pmatrix} A & B \\ C & D \end{pmatrix}$$

⇒

$$\Rightarrow \begin{pmatrix} [A, S^n], [B, S^n] \\ [C, S^n], [D, S^n] \end{pmatrix} = - \left[ \begin{pmatrix} A & B \\ C & D \end{pmatrix}, \frac{1}{2} z^n \right]$$

$$-\frac{1}{2} \left[ \begin{pmatrix} A & B \\ C & D \end{pmatrix}, \begin{pmatrix} 1 & \\ & -1 \end{pmatrix} \right] = \begin{pmatrix} 0 & B \\ C & 0 \end{pmatrix}, \quad -\frac{1}{2} \left[ \begin{pmatrix} A & B \\ C & D \end{pmatrix}, \begin{pmatrix} 1 & 1 \\ & 1 \end{pmatrix} \right] = -\frac{1}{2} \begin{pmatrix} B-C, A-D \\ D-A, C-B \end{pmatrix}$$

$$-\frac{1}{2} \left[ \begin{pmatrix} A & B \\ C & D \end{pmatrix}, \begin{pmatrix} i & \\ & -i \end{pmatrix} \right] = -\frac{1}{2} \begin{pmatrix} (B+C)i, -(A-D)i \\ (D-A)i, (B+C)i \end{pmatrix}$$

$$\Rightarrow [A+D, S^n] = 0, \text{ for } n=1, 2, 3$$

$$[B(u), S^3] = B(u), \quad [B(u), S^+] = [B(u), S^1] + i[B(u), S^2] \\ = -\frac{1}{2}(A-D) - \frac{1}{2}(A-D) = -(A-D)$$

or rewrite  $[S^n, A+D] = 0, [S^3, B(u)] = -B(u), [S^+, B(u)] = A(u) - D(u)$

Obviously  $S^+ |0\rangle = 0, S^3 |0\rangle = \frac{N}{2} |0\rangle$

$$S^3 B(u) = B(u)(S^3 - 1) \Rightarrow$$

$$S^3 B(u_1) \dots B(u_m) |0\rangle = B(u_1)(S^3 - 1) B(u_2) \dots |0\rangle$$

$$= B(u_1) B(u_2) (S^3 - 2) \dots |0\rangle = B(u_1) \dots B(u_m) (S^3 - M) |0\rangle$$

$$= \frac{N-2M}{2} B(u_1) \dots B(u_m) |0\rangle$$

$$[S^\dagger, B(v_1) \dots B(v_m)] = \sum_\alpha B(v_1) \dots [S^\dagger, B(u_\alpha)] \dots B(v_m)$$

$$= \sum_\alpha B(v_1) \dots [A(u_\alpha) - D(u_\alpha)] \cdot B(v_m)$$

$$\Rightarrow S^\dagger B(v_1) \dots B(v_m) |0\rangle = [S^\dagger, B(v_1) \dots B(v_m)] |0\rangle = \sum_\alpha B(v_1) \dots [A(u_\alpha) - D(u_\alpha)] \cdot B(v_m) |0\rangle$$

by using  $A(u) B(v) = \frac{1}{b(v-u)} B(v) A(u) - \frac{c(u-u)}{b(v-u)} B(u) A(v)$

$$D(u) B(v) = \frac{1}{b(u-v)} B(v) D(u) - \frac{c(u-v)}{b(u-v)} B(u) D(v)$$

$$D(u) |0\rangle = \prod_{j=1}^N b(u-u_j) |0\rangle, \quad A|0\rangle = |0\rangle.$$

We can expect that, the final expression can be expressed as

$$S^\dagger B(v_1) \dots B(v_m) |0\rangle = \sum_\alpha M_\alpha B(v_1) \dots B(v_{\alpha-1}) B(v_{\alpha+1}) \dots B(v_m) |0\rangle$$

The coefficient of  $M_1$  can be obtained as the term missing

$$M_1 \Rightarrow \prod_{\alpha=2}^m \frac{1}{b(v_\alpha - v_1)} - \frac{\prod_{j=1}^N b(v_1 - u_j)}{\prod_{\alpha=2}^m b(v_1 - v_\alpha)} = 0.$$

$\boxed{B(v_1)}$

← Bethe ansatz Eq.

Because all B commute, we can do arbitrary permutation

of B. Thus from  $M_1 = 0$ , we obtain  $M_\alpha = 0$ .

$$\Rightarrow \boxed{S^\dagger B(v_1) \dots B(v_m) |0\rangle = 0}$$