

## 10 Work and energy

§: work done to move a charge: (against electro-static force

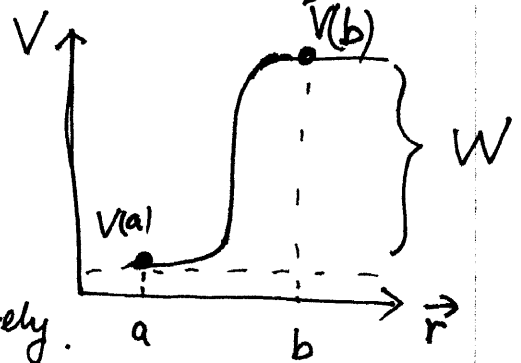
$$W = \int_a^b \vec{F} \cdot d\vec{l} = -Q \int_a^b \vec{E} \cdot d\vec{l} = Q [V(b) - V(a)]$$

A natural question is that:

Consider a charge distribution

$q_1, \dots, q_n$  located at  $r_1, \dots, r_n$ , respectively.

What's the energy stored in such a system?



This energy equals <sup>to</sup> the work done

during the process

the charges

$q_1, r_1$

$q_2, r_2$

of bringing

from well-separated <sup>long</sup> distance to the current distribution.

For example, let us consider two charges  $q_1$  and  $q_2$  at  $r_1$  and  $r_2$ .

The work to put them together: Let us first separate  $q_1, q_2$  at long distance. Then move  $q_1$  to  $r_1$ . No work is done during

this process. Then  $q_2$  is brought from infinity to  $r_2$ . The work

done

$$W_{12} = -q_2 \int_{+\infty}^{r_2} E d\vec{r} = +q_2 [V(r_2) - V(+\infty)] = \frac{q_2 q_1}{|r_1 - r_2|}$$

• if we have the third charge  $q_3$ , and bring it to  $r_3$ .

$$\begin{aligned} \text{the work done } W &= -q_3 \int_{-\infty}^{r_3} (\vec{E}_1 + \vec{E}_2) \cdot d\vec{\ell} = q_3 \left[ V_1(r_3) - V_1(\infty) \right. \\ &\quad \left. + V_2(r_3) - V_2(\infty) \right] \\ &= \frac{q_3 q_1}{|\vec{r}_3 - \vec{r}_1|} + \frac{q_3 q_2}{|\vec{r}_3 - \vec{r}_2|} \end{aligned}$$

add together, the energy of the three charge system

$$E = \sum W = \frac{q_1 q_2}{|\vec{r}_1 - \vec{r}_2|} + \frac{q_3 q_1}{|\vec{r}_3 - \vec{r}_1|} + \frac{q_3 q_2}{|\vec{r}_3 - \vec{r}_2|}$$

• please check that, if we have  $n$ -charge system  $(q_1, r_1), (q_2, r_2), \dots, (q_n, r_n)$ , then energy is

$$\begin{aligned} E &= \frac{q_1 q_2}{|\vec{r}_1 - \vec{r}_2|} \\ &+ q_3 \left( \frac{q_1}{|\vec{r}_3 - \vec{r}_1|} + \frac{q_2}{|\vec{r}_3 - \vec{r}_2|} \right) \\ &+ q_4 \left( \frac{q_1}{|\vec{r}_4 - \vec{r}_1|} + \frac{q_2}{|\vec{r}_4 - \vec{r}_2|} + \frac{q_3}{|\vec{r}_4 - \vec{r}_3|} \right) \\ &+ \dots \\ &+ q_n \left( \frac{q_1}{|\vec{r}_n - \vec{r}_1|} + \dots + \frac{q_{n-1}}{|\vec{r}_n - \vec{r}_{n-1}|} \right) \end{aligned} \left. \vphantom{\begin{aligned} E &= \dots \end{aligned}} \right\} \begin{aligned} &= \sum_{i < j} \frac{q_i q_j}{|\vec{r}_i - \vec{r}_j|} \\ &= \frac{1}{2} \sum_{i \neq j} \frac{q_i q_j}{|\vec{r}_i - \vec{r}_j|} \end{aligned}$$

avoid double counting.

$$\begin{aligned} \text{or } E &= \frac{1}{2} \sum_i q_i \left( \sum_{j \neq i} \frac{q_j}{|\vec{r}_j - \vec{r}_i|} \right) \\ &= \frac{1}{2} \sum_i q_i V(r_i) \end{aligned}$$

$V(r_i)$  is the potential at  $r_i$  generated from all other charges. (not include itself!)

Let consider a continuous <sup>charge</sup> distribution  $\rho(r)$

$$W = \frac{1}{2} \int \rho(r) V(r) d^3 r,$$

$$= \frac{1}{2} \int \frac{\rho(r) \rho(r')}{|\vec{r} - \vec{r}'|} d^3 \vec{r} d^3 \vec{r}'$$

for a continuous charge distribution, if  $V(r)$  is regular, we don't need to impose  $\vec{r}_i \neq \vec{r}'_i$ , as long as the integral converge.

§ Where is the energy stored? — charges or fields

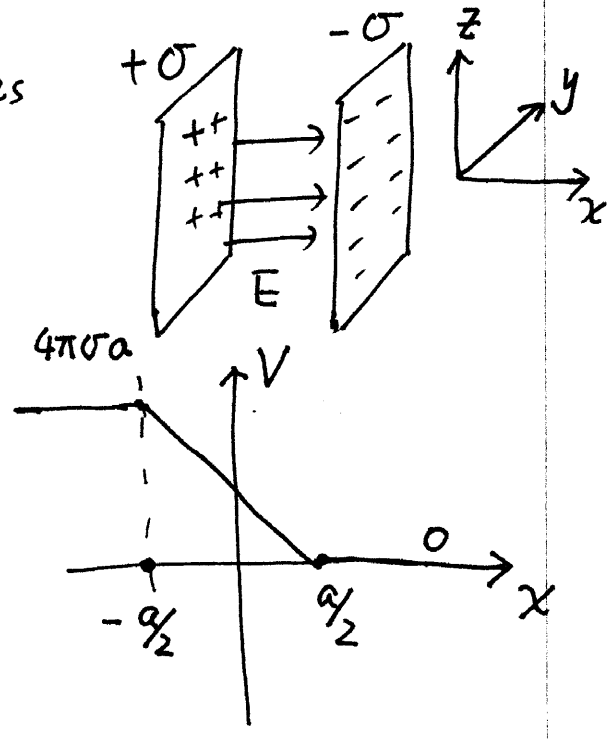
Let us consider two plates with opposite charge density

the E field is nonzero between two plates

$$\vec{E} = \begin{cases} 4\pi\sigma \hat{x}, & (-\frac{a}{2} < x < \frac{a}{2}) \\ 0 & \text{elsewhere} \end{cases}$$

⇒ the electric potential

$$V(x) = \begin{cases} 0, & x > \frac{a}{2} \\ 2\pi\sigma(a - 2x) & -\frac{a}{2} < x < \frac{a}{2} \\ 4\pi\sigma a & x < -\frac{a}{2} \end{cases}$$



$$\Rightarrow W = \frac{1}{2} \int \rho(r) V(r) d^3 r$$

$$\rho(r) = \sigma [\delta(x+a) - \delta(x-a)]$$

$$\Rightarrow W = \frac{\sigma}{2} \int dy dz \int_{dx} [\delta(x+a) - \delta(x-a)] V(x, y, z)$$

$$W = \frac{\sigma}{2} \int dydz \left[ V\left(\frac{a}{2}\right) - V\left(\frac{a}{2}\right) \right] = \frac{\sigma}{2} \cdot A \cdot 4\pi\sigma a = 2\pi\sigma^2 \text{Vol}$$

$$\Rightarrow \boxed{\frac{W}{\text{Vol}} = 2\pi\sigma^2 = \frac{E^2}{8\pi}}$$

Suppose that, we enlarge the distance between two plates, What's  
 the work we do? The energy density remains  $2\pi\sigma^2$ ,  
 but the volume changes  $A da$ ,  $\Rightarrow$  the total energy increase!

But what really changed during the process? The charge distribution on two plates doesn't change. It's much more naturally to attribute the energy change is due to the increase of the volume with electric field  $\vec{E}$ . In other words,

energy is stored in the fields, not just at the location of charges!

We need to change our ideology from charge to field.

Charge generates the field, but field gains its own life!

A formal derivation

$$W = \frac{1}{2} \int \rho(r) V(r) d^3r, \quad \leftarrow \quad \nabla \cdot \vec{E} = 4\pi\rho$$

$$= \frac{1}{8\pi} \int \nabla \cdot \vec{E} V(r) d^3r \quad \leftarrow \text{partial integral}$$

$$= \frac{1}{8\pi} \left[ \int \nabla \cdot (\vec{E} V(r)) d^3r - \int \vec{E} \cdot \nabla V(r) d^3r \right]$$

$$= \frac{1}{8\pi} \left[ \int \vec{E} \cdot \vec{E} d^3r + \oint_S V \vec{E}(r) \cdot d\vec{a} \right]$$

if we integrate over all the space, and set the radius of

the surface  $\rightarrow +\infty$ . We know  $V \rightarrow \frac{1}{R}$ ,  $E \rightarrow \frac{1}{R^2}$

the area  $\propto R^2$

$\Rightarrow$  the last term  $\propto \frac{1}{R} \rightarrow 0$ .

$$\Rightarrow W = \frac{1}{8\pi} \int E^2 d^3r$$

or energy density

$$\boxed{\rho_w = \frac{E^2}{8\pi}}$$

Remarks

① What's the energy stored in a point, say, electron?  
the field

$$E = \frac{e}{r^2}$$

$$\Rightarrow W = \frac{1}{8\pi} \int_{r_0}^{+\infty} \frac{e^2}{r^4} \cdot 4\pi r^2 dr$$

$$= \frac{e^2}{8\pi} \cdot 4\pi \int_{r_0}^{+\infty} \frac{dr}{r^2} = \frac{e^2}{2r_0} \rightarrow +\infty \text{ as } r_0 \rightarrow 0.$$

This is indeed a difficulty. Actually, in order to confine charge in an infinitesimal volume, this kind of divergence is unavoidable.

Actually, there's a way to estimate electron radius as

$$\frac{e^2}{r_0} = m_e c^2, \quad r_0 = \frac{m_e c^2}{e^2}, \text{ which is called electron } \text{em-radius.}$$

$$\sim 2.8 \times 10^{-15} \text{ m}$$

But actually, we know

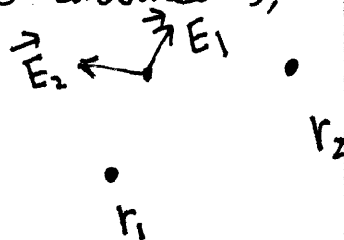
electron size is much smaller than  $r_0$ , if it's not zero.

Nevertheless, this kind of energy divergence is a frozen of point charge energy. It does n't change, during

any process. We only care about the energy difference,

thus this is not a problem to us.

If we approximate ~~an~~ charge distribution as continuous, this difficulty is avoid.



For two charges, the energy stored

$$W_{12} = \frac{1}{8\pi} \left[ \int (\vec{E}_1 + \vec{E}_2)^2 dr - \int \vec{E}_1^2 dr - \int \vec{E}_2^2 dr \right]$$

$$W_{12} = \frac{1}{4\pi} \int \vec{E}_1 \cdot \vec{E}_2 d^3r$$

The energy associated with the point charge  $q_1, q_2$  are subtracted  
 divergent

check  $W_{12} = \frac{-1}{4\pi} \int \vec{E}_1 \cdot \nabla V_2(r) d^3r$

$$= \frac{-1}{4\pi} \int d^3r \nabla (\vec{E}_1 \cdot V_2(r)) - V_2(r) \nabla \cdot \vec{E}_1 d^3r$$

$$= \frac{-1}{4\pi} \left[ \oint d\vec{a} \cdot \nabla (\vec{E}_1(r) \frac{q_2}{|\vec{r}-\vec{r}_2|}) \right] + \frac{1}{4\pi} \int d^3r \frac{q_2}{|\vec{r}-\vec{r}_2|} \frac{q_1}{4\pi |\vec{r}-\vec{r}_1|}$$

$\swarrow$   
 $= 0$

$\downarrow$   
 $\frac{q_1 q_2}{|\vec{r}_1 - \vec{r}_2|}$

set surface to infinity

$$\left. \begin{matrix} E \rightarrow \frac{1}{R^2} \\ V \rightarrow \frac{1}{R} \\ \oint da \rightarrow R^2 \end{matrix} \right\} \frac{1}{R} \rightarrow 0$$

$\Rightarrow W_{12} = \frac{q_1 q_2}{|\vec{r}_1 - \vec{r}_2|}$  ✓