

Lect 7: - Nested Bethe Ansatz
Spin-1/2 fermion

The model has two conserved quantities: the total particle number N and the total spin S_z . Correspondingly, we have two sets of BA quantum numbers and two sets of BA equations

$$\textcircled{1} \quad e^{ik_j L} = \prod_{\alpha=1}^M \frac{k_j - \Lambda_{\alpha} + ic/2}{k_j - \Lambda_{\alpha} - ic/2} \quad \text{for each } k_j, j=1, 2, \dots, N$$

(particle number)

$$\textcircled{2} \quad \prod_{j=1}^N \frac{k_j - \Lambda_{\alpha} - ic/2}{k_j - \Lambda_{\alpha} + ic/2} = - \prod_{\beta=1}^M \frac{\Lambda_{\beta} - \Lambda_{\alpha} - ic}{\Lambda_{\beta} - \Lambda_{\alpha} + ic} \quad \alpha=1, \dots, M$$

(# of spin up)

k_j ($j=1, 2, \dots, N$) and Λ_{α} ($\alpha=1, \dots, M$) are two sets of physical quantities that we are going to solve.

Eq ① \Rightarrow

$$k_j L = 2\pi \cdot \text{Integer} + \sum_{\alpha=1}^M 2 \tan^{-1} \frac{c}{2(k_j - \Lambda_{\alpha})}$$

$$= 2\pi \cdot \text{integer} + \sum_{\alpha=1}^M \pi - 2 \tan^{-1} \frac{2(k_j - \Lambda_{\alpha})}{c}$$

$$k_j L = 2\pi I_j - 2 \sum_{\alpha=1}^M \tan^{-1} \frac{2(k_j - \Lambda_{\alpha})}{c} \quad (j=1, 2, \dots, N)$$

I_j : integer if M is even;
half integer if M is odd.

Eq ② \Rightarrow

$$\sum_{j=1}^N 2 \tan^{-1} \frac{c}{2(k_j - \Lambda_{\alpha})} = 2\pi \cdot \text{integer} + \pi + \sum_{\beta=1}^M 2 \tan^{-1} \frac{c}{\Lambda_{\beta} - \Lambda_{\alpha}}$$

$$\sum_{j=1}^N \pi - \sum_{j=1}^N 2 \tan^{-1} \frac{2(k_j - \Lambda_{\alpha})}{c} = 2\pi [\text{half integer}] + \sum_{\beta=1}^M \pi - \sum_{\beta=1}^M 2 \tan^{-1} \frac{\Lambda_{\beta} - \Lambda_{\alpha}}{c}$$

$$- \sum_{j=1}^N 2 \tan^{-1} \frac{2(k_j - \Lambda_\alpha)}{c} = 2\pi \left(\text{half integer} + \frac{N-M}{2} \right) - \sum_{\beta=1}^M 2 \tan^{-1} \frac{\Lambda_\beta - \Lambda_\alpha}{c}$$

$$2 \sum_{j=1}^N \tan^{-1} \frac{2(k_j - \Lambda_\alpha)}{c} = -2\pi J_\alpha + 2 \sum_{\beta=1}^M \tan^{-1} \frac{\Lambda_\beta - \Lambda_\alpha}{c} \quad \alpha = 1, \dots, M$$

where $J_\alpha = \begin{cases} \text{integer, if } N-M \text{ is odd} \\ \text{half-integer if } N-M \text{ is even} \end{cases}$

Ground state solution:

① First consider the limit $c \rightarrow \infty$, $N = \text{even}$, $m = \text{odd}$

$$a) \quad k_j L = 2\pi I_j + 2 \sum_{\beta=1}^M \tan^{-1} \frac{2\Lambda_\beta}{c} \quad j = 1, \dots, N$$

$$b) \quad -2N \tan^{-1} \frac{2\Lambda_\alpha}{c} = 2\pi J_\alpha + 2 \sum_{\beta=1}^M \tan^{-1} \frac{\Lambda_\beta - \Lambda_\alpha}{c} \quad \alpha = 1, \dots, M$$

I_j : half integer. J_α : integer

$$b) \Rightarrow -2N \sum_{\alpha=1}^M \tan^{-1} \frac{2\Lambda_\alpha}{c} = -2\pi \sum_{\alpha=1}^M J_\alpha \quad \leftarrow \sum_{\alpha, \beta} \tan^{-1} \frac{\Lambda_\beta - \Lambda_\alpha}{c} = 0$$

$$\text{hence} \quad 2 \sum_{\alpha=1}^M \tan^{-1} \frac{2\Lambda_\alpha}{c} = \frac{2\pi}{N} \sum_{\alpha=1}^M J_\alpha$$

$$k_j = \frac{2\pi}{L} (I_j + a), \quad \text{with } a = -\frac{1}{N} \sum_{\alpha=1}^M J_\alpha$$

$$E = \sum_{j=1}^N k_j^2 = \left(\frac{2\pi}{L} \right)^2 \sum_{j=1}^N (I_j + a)^2$$

Ground state: when $a=0$, and I_j are closely packed around 0 symmetrically.

* $a=0$ can be viewed as the center of mass momentum.

$(I_j + a)^2 \propto$ the i th particle kinetic energy.

$a=0 \Rightarrow J_\alpha$ takes value symmetrically with respect to zero ③

from b) $\Rightarrow \Lambda_\alpha$ and $-\Lambda_\alpha$ appear in pairs.

2) $c \neq \infty$. in the ground state J_j still takes the value as the case of $c = \infty$. J_α distributes symmetrically around the origin, and Λ and $-\Lambda$ paired.

$$2\pi J_j = L k_j + \sum_\alpha 2 \tan^{-1} \frac{2k_j - 2\Lambda_\alpha}{c} = L k_j + \sum_\alpha \tan^{-1} \frac{2k_j - 2\Lambda_\alpha}{c} + \tan^{-1} \frac{2k_j + 2\Lambda_\alpha}{c}$$

$$2\pi J_j = L k_j + \sum_\alpha \tan^{-1} \frac{4k_j c}{c^2 + 4\Lambda_\alpha^2}$$

where Λ_α^2 increases, k_j increases \Rightarrow where Λ_α closely packed around zero, we arrive at the ground state

we consider the case $N, L, M \rightarrow \infty$, but $N/L, M/L$ fixed.
the ground state in

$$J_{j+1} - J_j = 1, \quad J_{\alpha+1} - J_\alpha = 1$$

define

$$L \rho(k) dk = \# \text{ of } k_j \text{ in the interval } k \rightarrow k+dk$$

$$L \sigma(\Lambda) d\Lambda = \# \text{ of } \Lambda_\alpha \text{ in the interval } \Lambda \rightarrow \Lambda + d\Lambda.$$

$$f = I/L \quad \text{and} \quad g = J/L$$

$$\Rightarrow \frac{df}{dk} = \rho(k) \quad \text{and} \quad \frac{dg}{d\Lambda} = \sigma(\Lambda)$$

$$\text{From } k_j L = 2\pi J_j + \sum_{\alpha=1}^M \theta(2k_j - 2\Lambda) \quad \text{where } \theta(x) = -2 \tan^{-1} x/c$$

$$\rightarrow k = 2\pi f + \int_{-B}^B \theta(2k - 2\Lambda) \sigma(\Lambda) d\Lambda$$

$$\frac{d}{dk} \rightarrow \boxed{2\pi p(k) = 1 + \int_{-B}^B \frac{4c \sigma(\Lambda)}{c^2 + 4(k-\Lambda)^2} d\Lambda}$$

by using $\frac{d\theta}{dx} = -2 \frac{c}{(c^2+x^2)}$

$$2 \sum_{j=1}^N \tan^{-1} \frac{2(k_j - \Lambda_\alpha)}{c} = -2\pi J_\alpha + 2 \sum_{\beta=1}^M \tan^{-1} \frac{\Lambda_\beta - \Lambda_\alpha}{c}$$

$$\sum_{j=1}^N \theta(2\Lambda_\alpha - 2k_j) = -2\pi J_\alpha + \sum_{\beta=1}^M \theta(\Lambda_\alpha - \Lambda_\beta)$$

$$\int_{-Q}^Q \theta(2\Lambda - 2k) p(k) dk = -2\pi g + \int_{-B}^B \theta(\Lambda - \Lambda') \sigma(\Lambda') d\Lambda'$$

$$\frac{d}{d\Lambda} \Rightarrow \boxed{2\pi \sigma(\Lambda) = - \int_{-B}^B \frac{2c \sigma(\Lambda')}{c^2 + (\Lambda - \Lambda')^2} d\Lambda' + \int_{-Q}^Q \frac{4c p(k) dk}{c^2 + 4(k - \Lambda)^2}}$$

Also the constraints

$$\boxed{\frac{N}{L} = \int_{-Q}^Q p(k) dk, \quad \frac{M}{L} = \int_{-B}^B \sigma(\Lambda) d\Lambda}$$

we can solve p, σ, Q, B , then

$$\frac{E}{L} = \int_{-Q}^Q k^2 p(k) dk$$